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# LEM/ODT Multiphase Applications

**Alan Kerstein**  
**Consultant**  
**Danville CA, USA**

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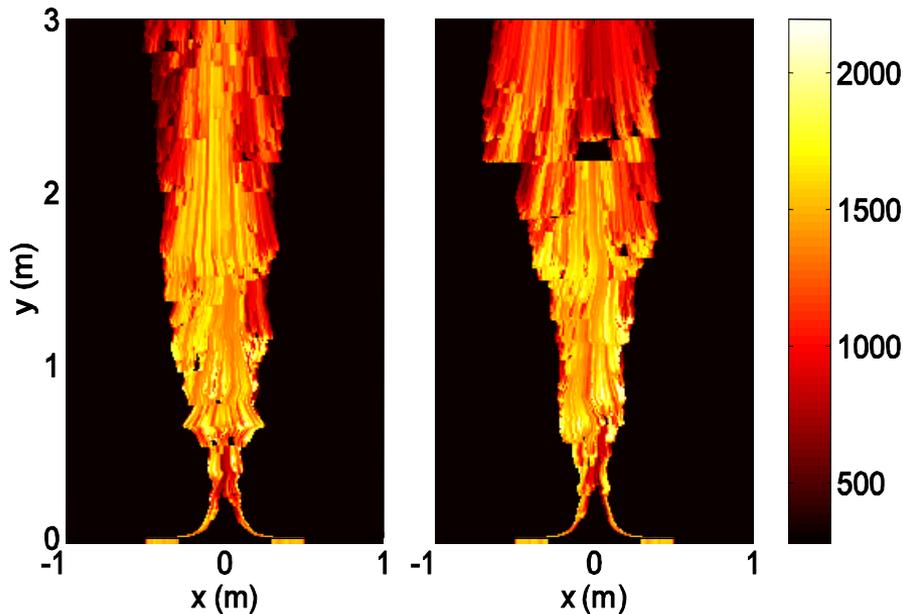
**MetStröm Short Course at BTU**

# Outline of presentation

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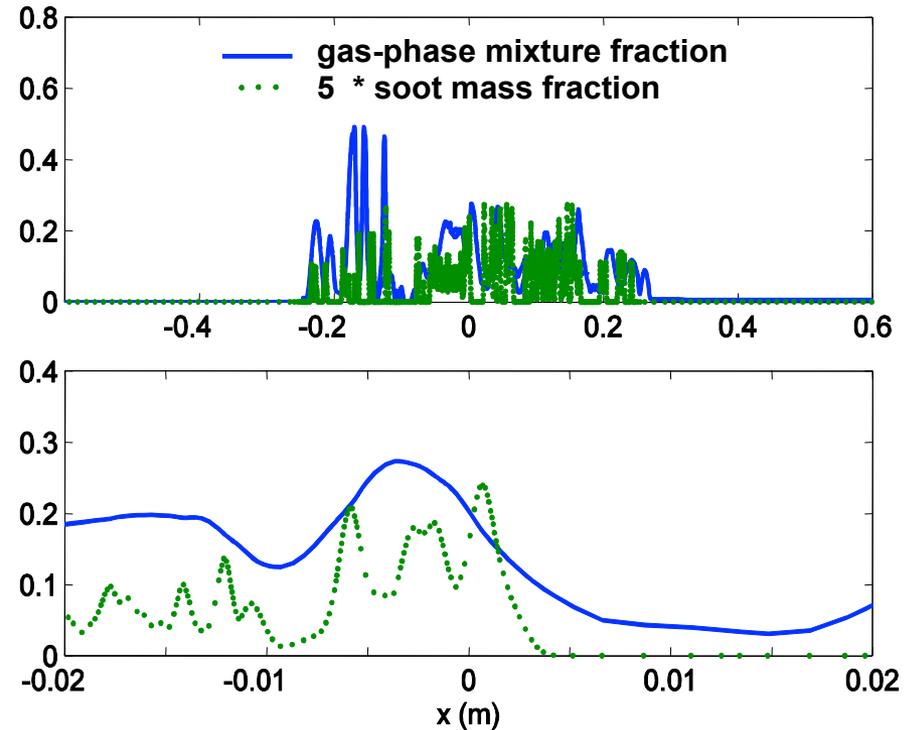
- Sooting plume
- Wall deposition
- Clustering

# Adaptive-mesh ODT was used to simulate an ethylene-air sooting plume



## An effect captured by spatial advancement:

The spatial continuity equation induces narrowing of temperature fields above the inlet due to lateral inflow balancing vertical buoyant acceleration



The adaptive mesh efficiently resolves small features – the new code allows different meshes for different properties, e.g., high-Sc scalars, enabling a big time-step increase

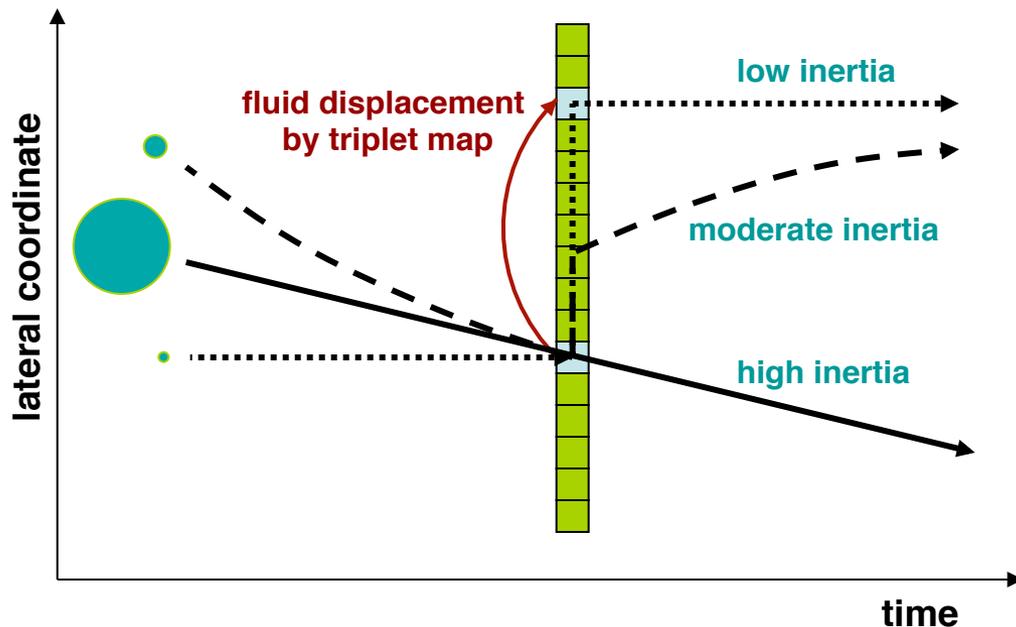
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# A particle-eddy interaction couples entrained particles to fluid motion (one-way coupling)

- In ODT, motion and velocity are distinct, though dynamically consistent
- Particles respond, via drag law, to motion (in ODT, eddy events)
- Because ODT eddies are instantaneous
  - an internal (eddy) time coordinate for particle-eddy interaction is introduced
  - this involves another free parameter, relating the interaction time to  $t$

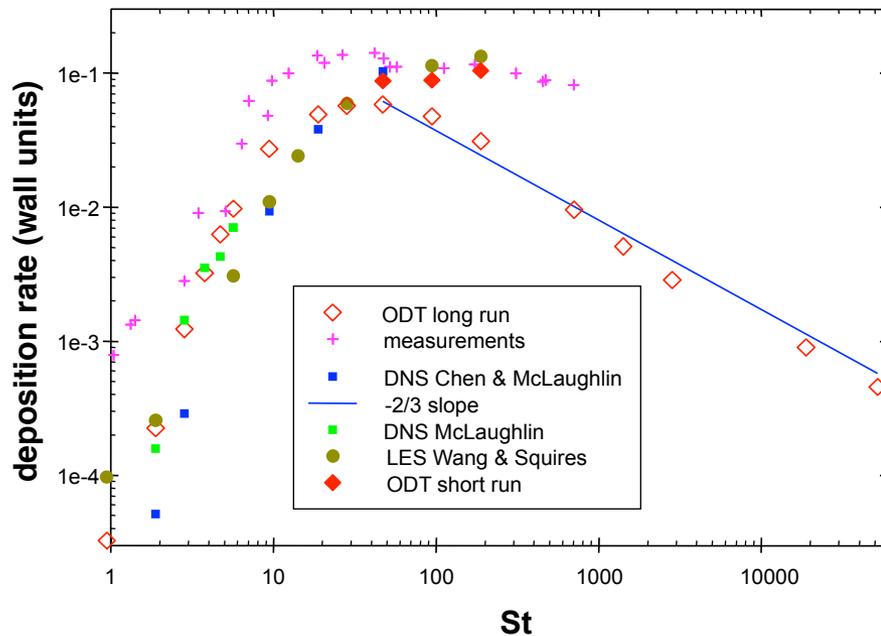


- **Eddy-time integration determines a trajectory 'jump condition' representing the eddy-induced trajectory change, adjusted so future motion is not double-counted**
- **Ballistic motion remains linear**
- **Zero-inertia (no-slip) particles follow the fluid**
- **Particle-fluid relative motion is realistic, though absolute motion is discontinuous**

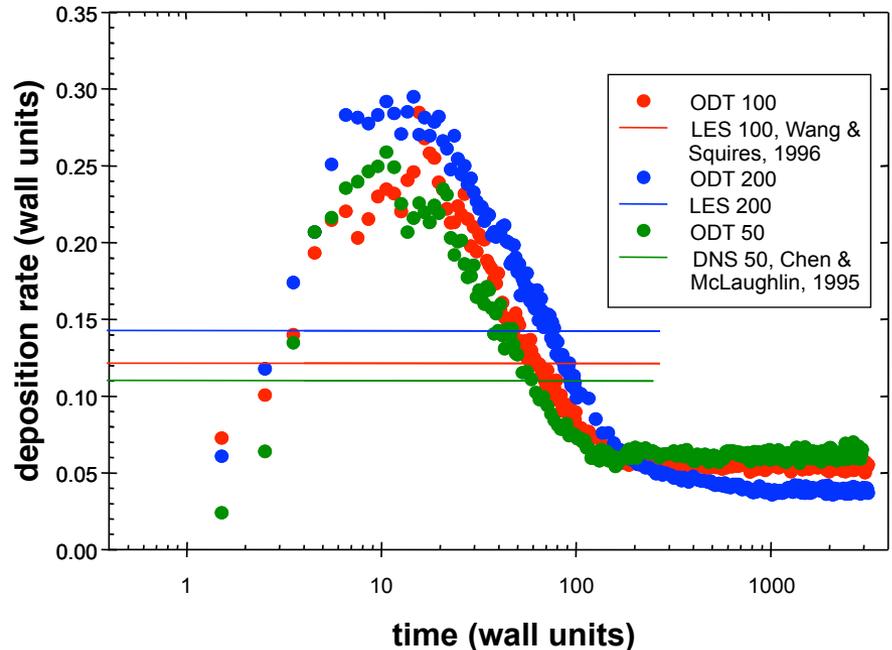
# Measured and 3D-simulated wall deposition is reproduced, and a new regime is found

Wall deposition in *turbulent channel flow* (the ODT domain is wall-normal)

Dependence on Stokes number

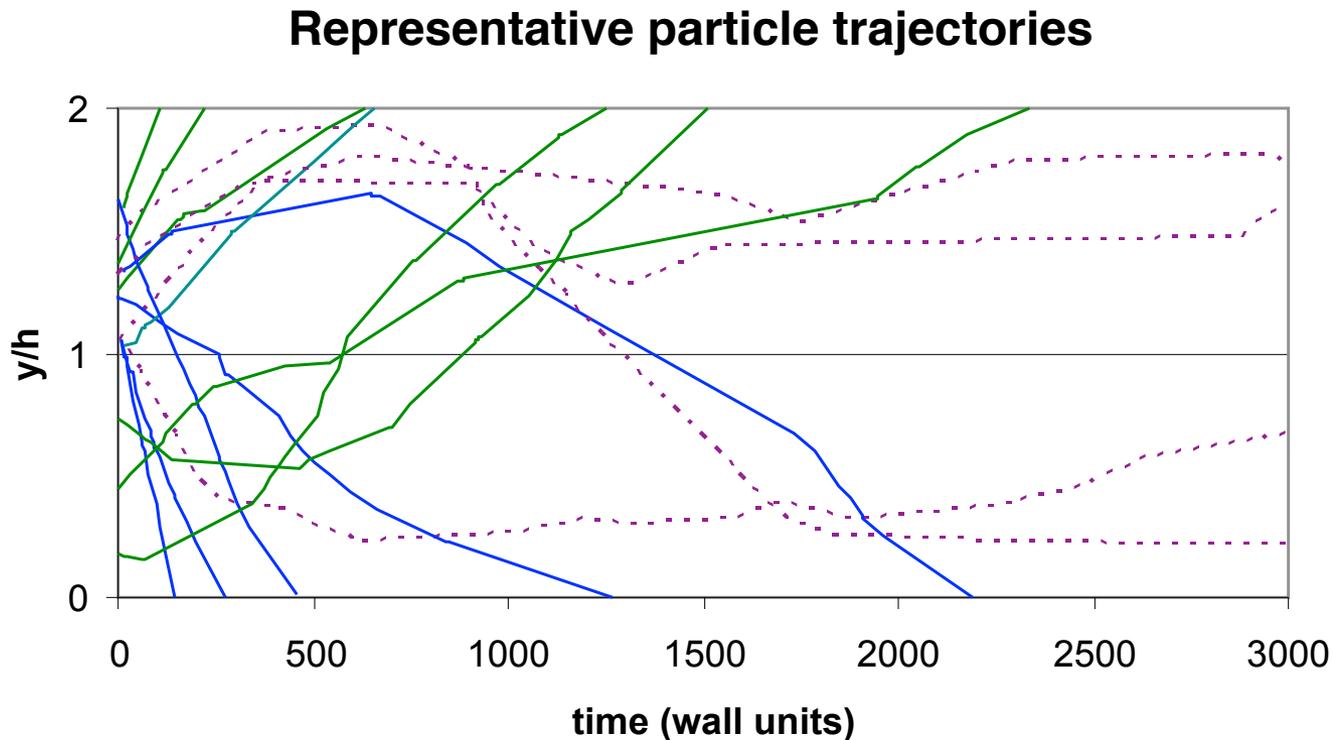


Time variation of deposition rate (transient relaxation)



**Comparisons suggest that measurements and 3D simulations are seeing initial transients rather than the late-time regime indicated by ODT**

# Early deposition is ballistic, late deposition is Stokes-number dependent



**The  $-2/3$  power dependence on  $St$  is explained by a simple scaling analysis. Closure analysis gives a much milder decline – and is ‘validated’ by data that mainly reflects initial conditions!**

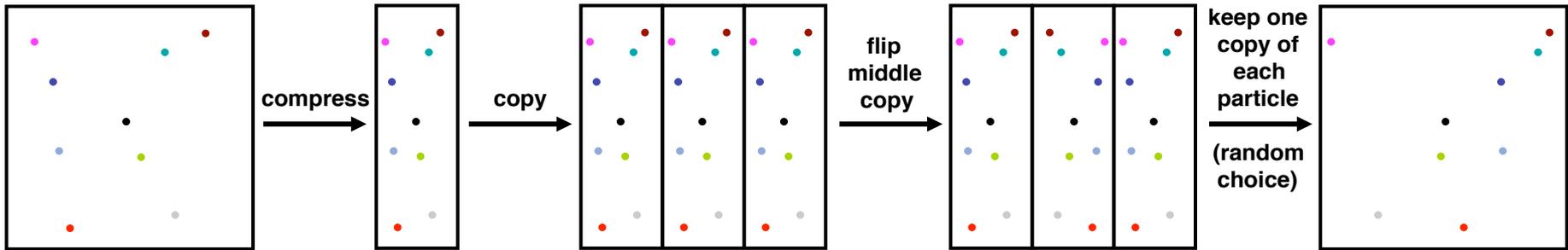
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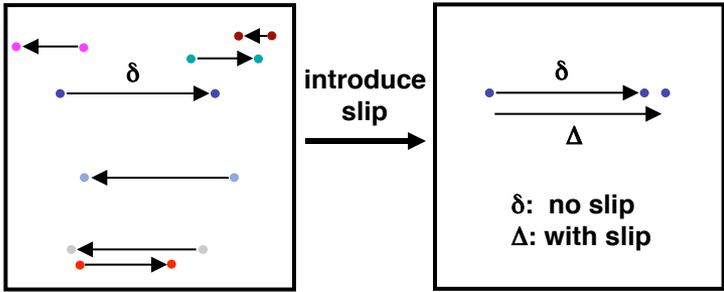
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# Using map-based advection, a 3D Lagrangian (grid-free) low-inertia particle advancement model is formulated

**Displacement of slip-free (zero-inertia) particles by a 3D triplet map:**



**Fluid displacements  $\delta$  are multiplicatively incremented to represent particle inertia:**



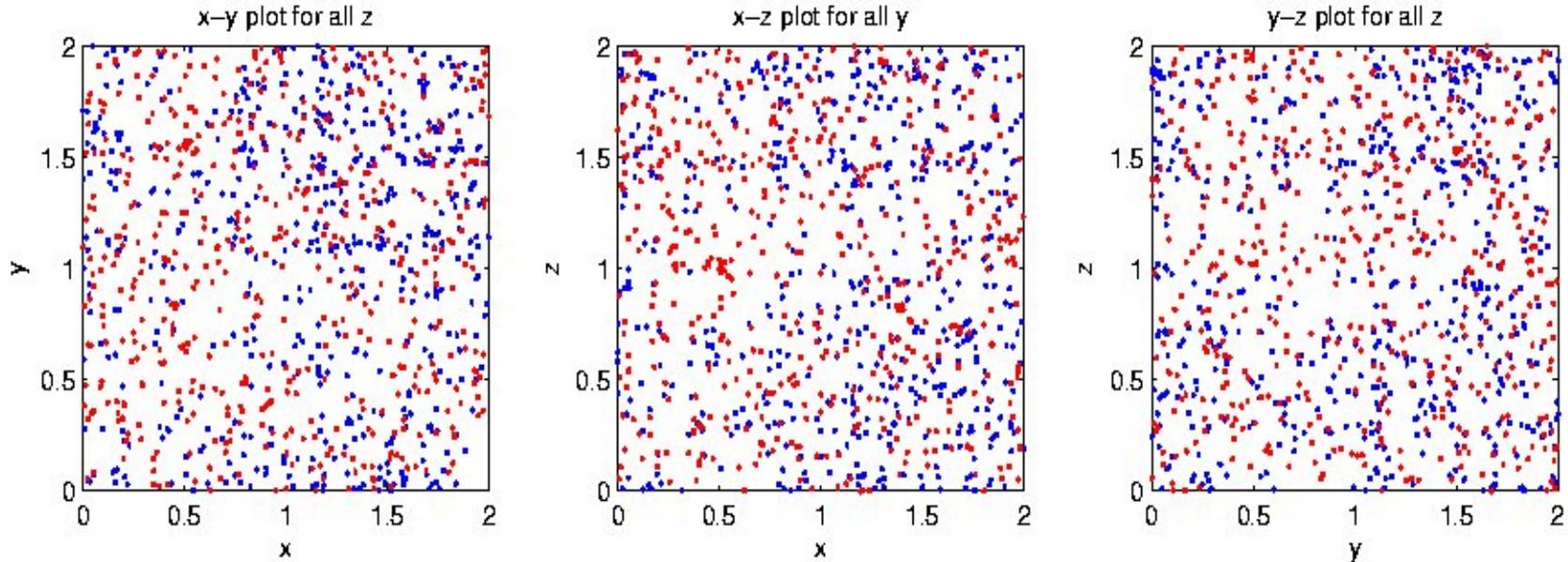
Inertia model:

$$\Delta = (1+S) \delta$$

$S \ll 1$  is the model analog of Stokes number,  
 $St = [\text{particle response time}] / [\text{flow time}]$

If polydisperse, S can be different for each particle

# For nonzero $S$ , clustering is observed

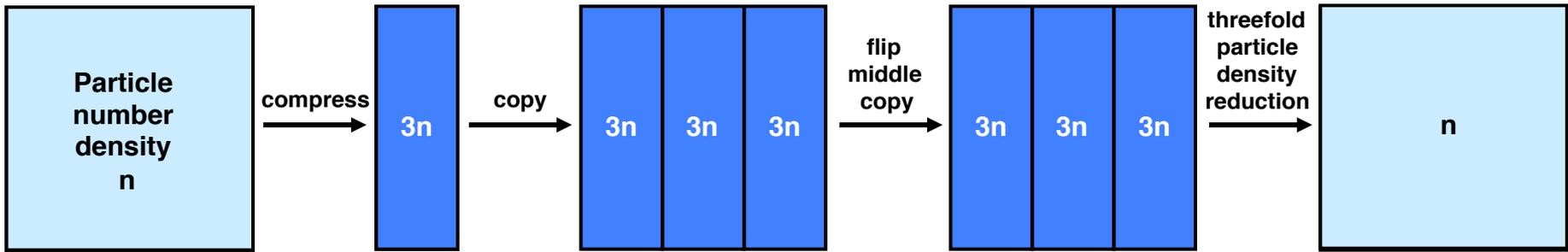


## Simulation:

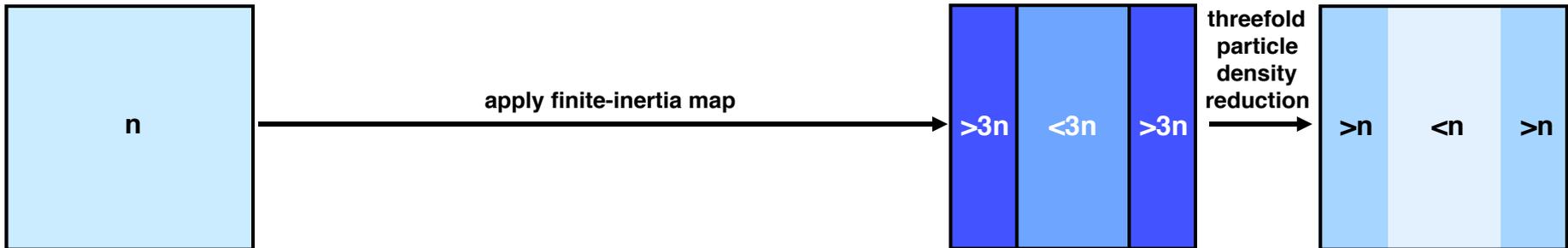
- Cubic domain, map size = domain size
- Maps in x, y and z directions, randomly positioned
- Periodic boundary conditions
- Iterated to statistical steady state
- Red,  $S = 0$ ; blue,  $S = 0.1$

# Continuum interpretation: slip induces fluctuations in an initially uniform particle-density field

**Zero inertia: uniform multiplicative compression, compensated by number reduction**



**Non-zero inertia: non-uniform compression, inducing particle-density fluctuations**



# Exact analysis yields parameter dependence of a clustering metric

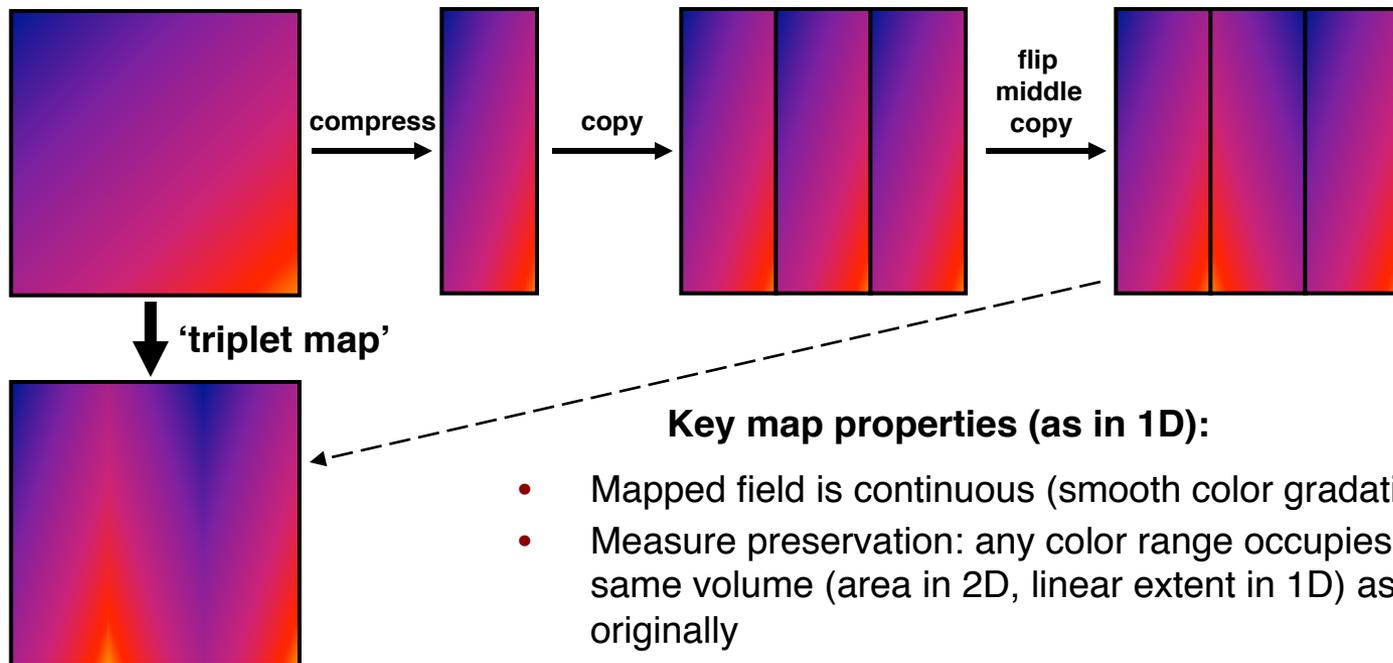
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- Radial distribution function (RDF)  $g(r)$ :
  - Likelihood of finding a particle at a distance  $r$  from a given particle
  - Normalized so  $g=1$  for statistically independent particles
- Prediction:
  - $g \sim r^{-cS_1S_2}$  for particles, labeled 1 and 2, with different  $S$  values
  - Valid for a restricted  $r$  range dependent on  $|S_1 - S_2|$  and flow structure
  - Previously obtained heuristically and with DNS (e.g., Chun et al. 2005)

# Significance (1): the analysis elucidates the geometrical basis of clustering

- Slip proportionality to displacement leads to the power-law  $r$  dependence of  $g$
- Clustering is a second-order effect (bilinear in  $S$ ) for continuous maps

Application of the advective map to an arbitrary continuous field:



# Significance (2): model properties suggest an efficient algorithm for simulation of particle motion

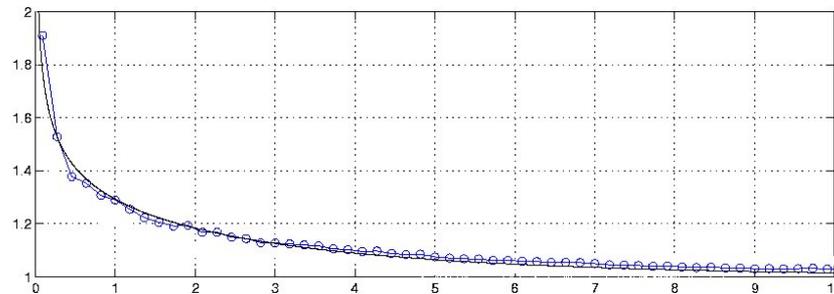
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- Motivation: turbulence enhancement of droplet coalescence
  - Collision rates are proportional to  $n^2$  locally, hence greater if  $n$  fluctuates
  - Gillespie's (1975) Stochastic Simulation Algorithm (SSA) captures collision randomness but not clustering effects
  - Map method captures both at no greater cost
- Application in progress (with Steve Krueger, U. of Utah): rain formation
  - Each raindrop that falls gathers a million others (snowball effect)
  - The one per million droplet that grows big enough to fall is rate controlling
  - Rare events (rapid coalescence) dominate, so need detailed simulations
- Future work: Embed this model in a simulation of larger-scale processes (explained shortly)

# Benchmarked the 3D model using DNS data, will imbed it in a multi-process cloud representation

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- Benchmarking:
  - Have tuned to match monodisperse (below) and bidisperse RDFs.
- Cloud application: simulate small scales in a 1D map-based scheme
  - Krueger's 1D EPM captures condensational growth in fluctuating humidity
  - Coalescence variability is important at smaller scales
  - Therefore structure the 1D scheme as a stack of cubes; 3D evolution in each
  - Sedimentation and droplet collision phenomenology have been incorporated

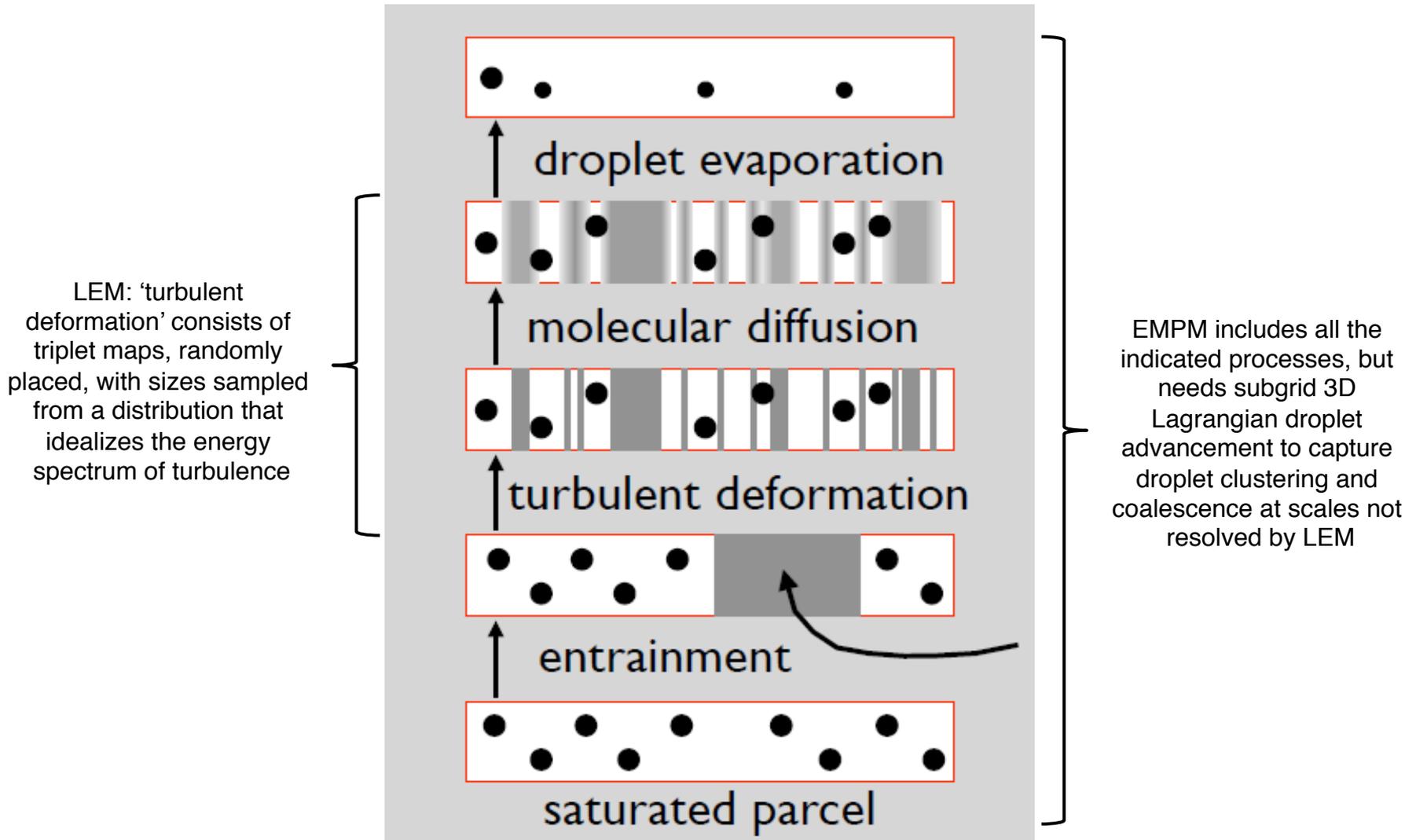


$g$  vs.  $r/[\text{Kolmogorov microscale}]$  for  $St=0.136$ .

Symbols, model; smooth curve, functional fit to DNS (Reade and Collins, 2000).

# The 1D Explicit Mixing Parcel Model (EMPM)

incorporates entrainment and phase change into LEM



# EMPM flow states resemble (and help interpret) measured data traces

## EMPM water vapor and temperature fields

